

Exact Diagonalization Bath Fitting

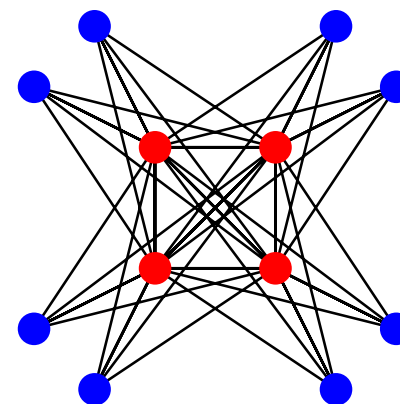
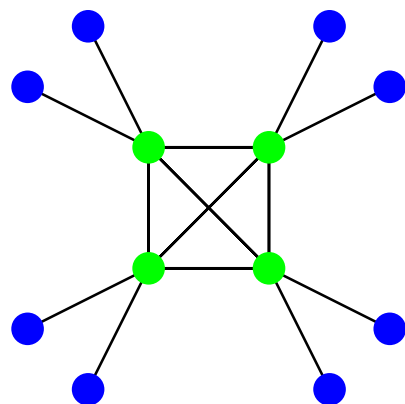
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Single Impurity Anderson Model

$$\mathcal{H} = \sum_{d,\sigma} \epsilon_d c_{d\sigma}^\dagger c_{d\sigma} + \frac{1}{2} \sum_{ijkl,\sigma\sigma'} U_{ijkl} c_{i\sigma}^\dagger c_{j\sigma'}^\dagger c_{l\sigma'} c_{k\sigma} +$$
$$\sum_{k,\sigma} \epsilon_k c_{k\sigma}^\dagger c_{k\sigma} + \sum_{kd,\sigma} \left[V_{kd} c_{k\sigma}^\dagger c_{d\sigma} + h.c. \right]$$

Single Impurity Anderson Model

- Diagonal basis
 - Each bath orbital only hybridizes with single impurity orbital
- Off-diagonal basis
 - All bath orbitals are coupled to every impurity orbital



*A. Liebsch and H. Ishida 2012 *J. Phys.: Condens. Matter* **24** 053201

$$G_{ii}^{imp,0}(\omega) = [(\omega + i\delta) - H^0 - \Delta_{ii}(\omega)]^{-1}$$

$$\Delta_{ij}(\omega) = \sum_k \frac{V_{ki} V_{kj}^*}{\omega + i\delta - \epsilon_k}$$

- Suitable for diagonal basis
- Two possible fitting options:
 - Green's function fitting (Density of States)
 - Direct fitting of Hybridization function (if available)
- Broadening improves fitting

- Green's function fitting

$$\left| \operatorname{Im} G_{ii}^{imp,0}(\omega) - \operatorname{Im} \left\{ \frac{1}{\omega + i\delta - \epsilon_i - \sum_{k=1}^{N_s} \frac{V_{ki}^2}{\omega + i\delta - \epsilon_k^i}} \right\} \right| \rightarrow \min$$
$$\{\epsilon_i, \epsilon_k^i, V_{ik}\}$$

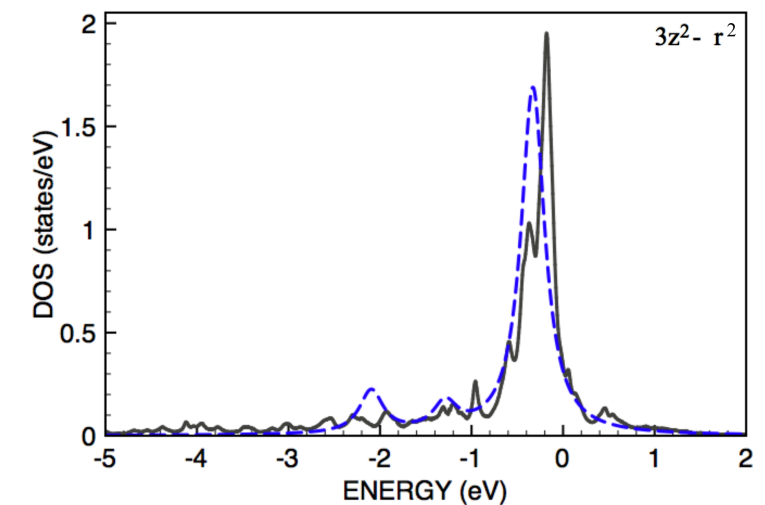
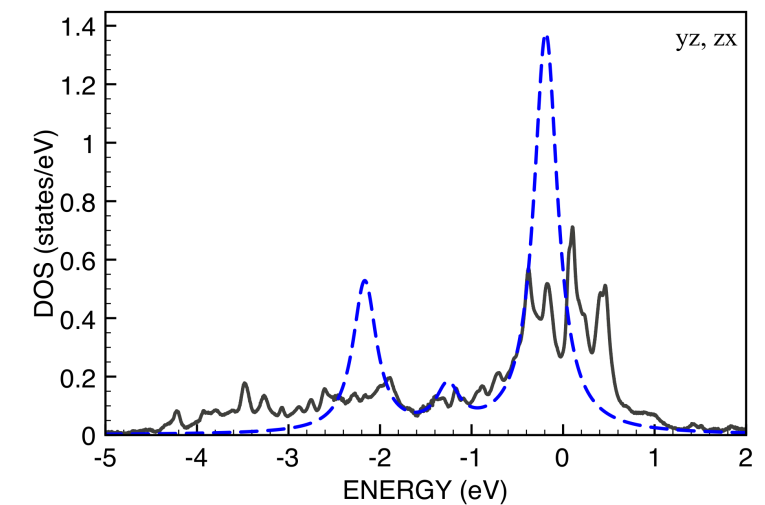
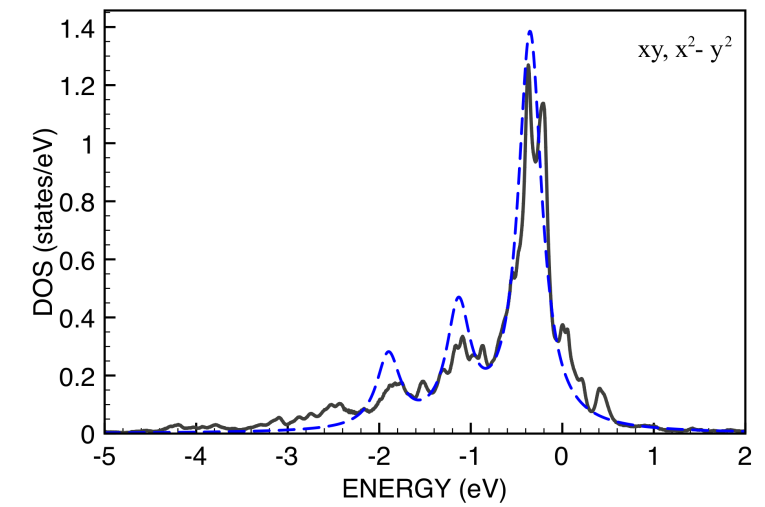
- Hybridization function fitting

$$\left| \operatorname{Im} \Delta_{ii}(\omega) - \operatorname{Im} \left\{ \sum_{k=1}^{N_s} \frac{V_{ki}^2}{\omega + i\delta - \epsilon_k^i} \right\} \right| \rightarrow \min$$
$$\{\epsilon_k^i, V_{ik}\}$$

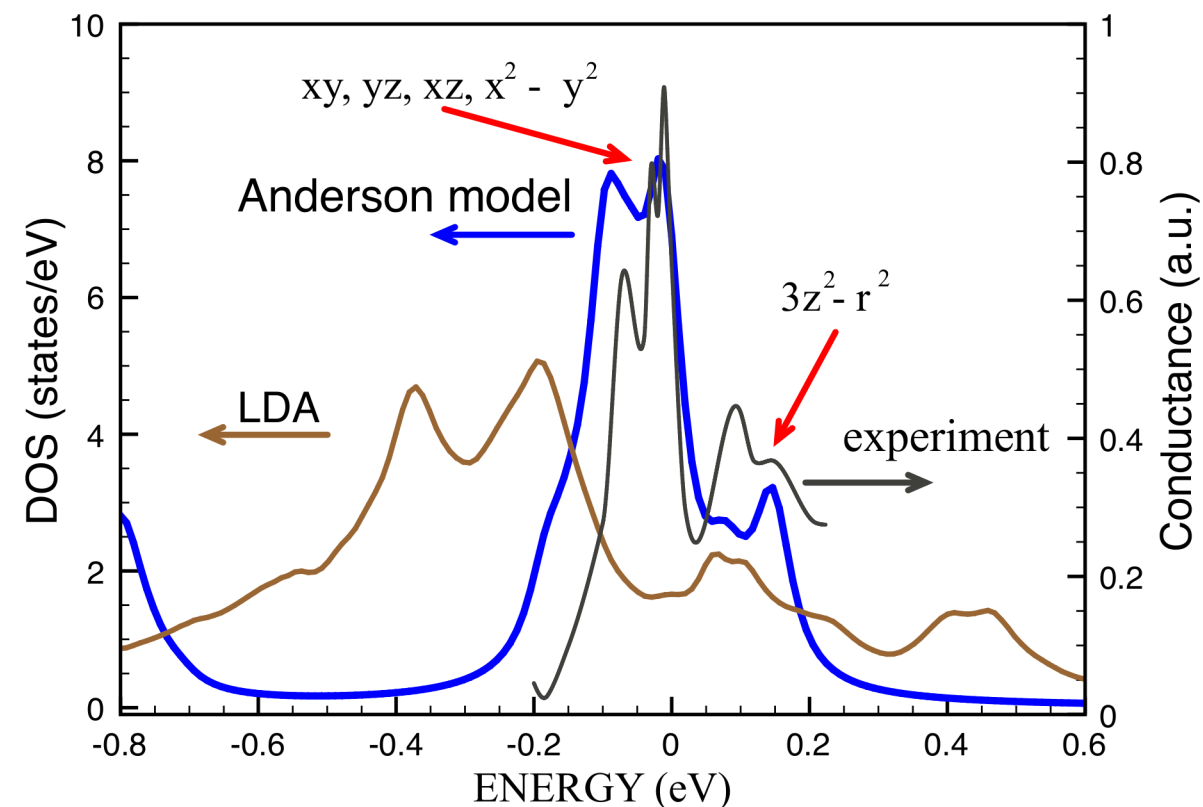
- ϵ_i are chosen to preserve target density

Real Frequency fitting

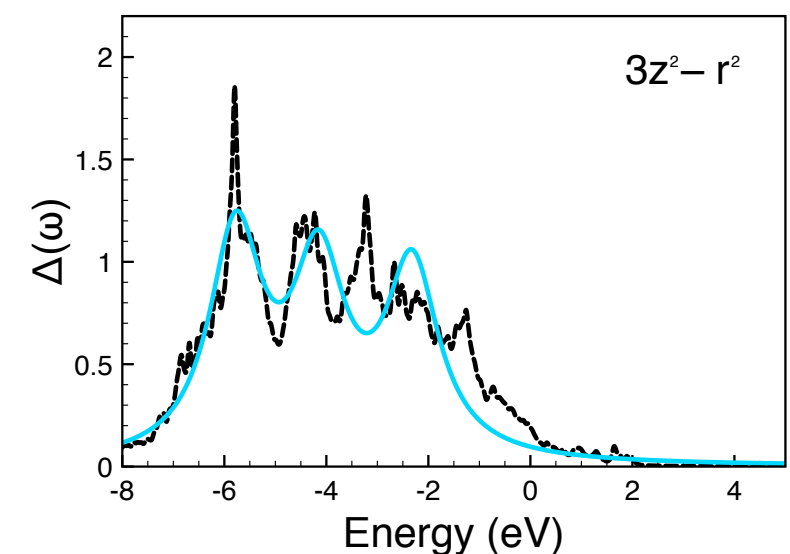
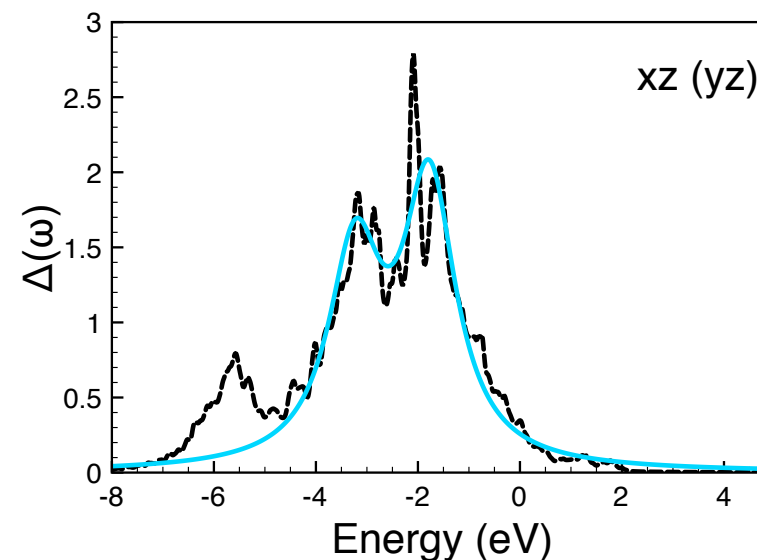
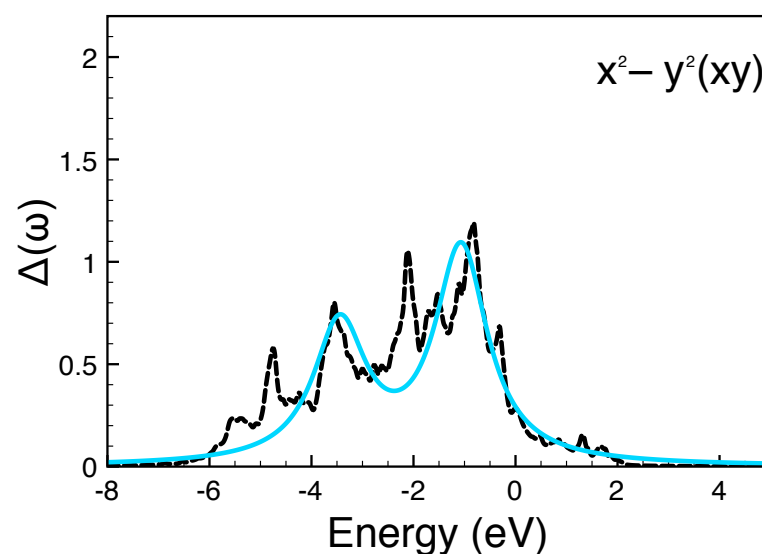
- Green's function fitting
 - Can capture main features of non-interacting density of states
 - Can reproduce main features of interacting spectra



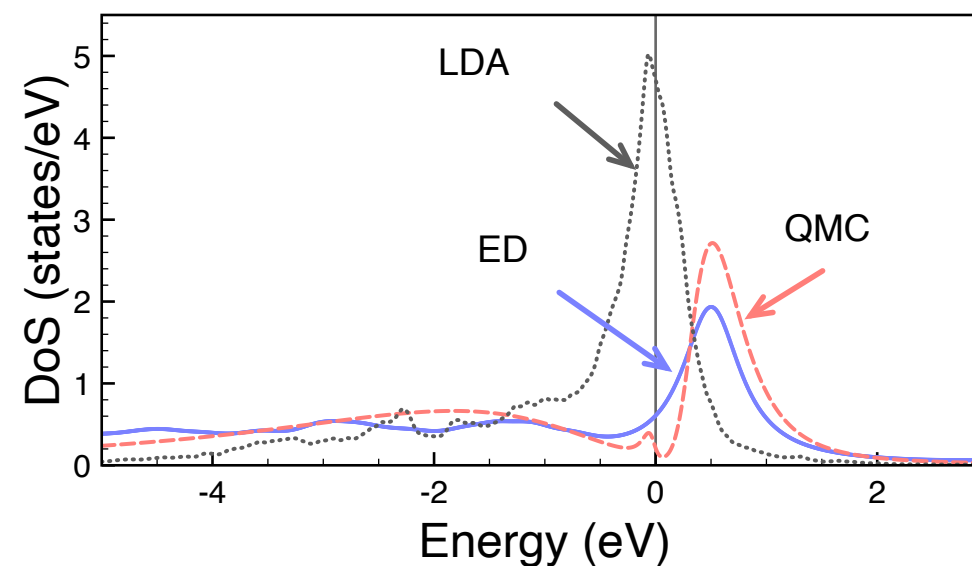
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- Hybridization function fitting
 - Allow better control of low-energy features
 - Orbital energies are chosen to maintain target density



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- Pros
 - Provides control to features that need to be captured
 - Allows to work directly with DFT data
- Cons
 - Fitting result strongly depends on initial parameters
 - Extremely hard for automatic fitting

- Data is available on Matsubara axis (DMFT or SEET)

$$\mathbf{G}(i\omega_n) = \left[(i\omega_n) - H^0 - \Delta(i\omega_n) \right]^{-1}$$

$$\Delta_{ij}(i\omega_n) = \sum_k \frac{V_{ik} V_{jk}^*}{i\omega_n - \epsilon_k} \approx \sum_{k=1}^{N_s} \frac{V_{ik} V_{jk}^*}{i\omega_n - \epsilon_k}$$

- Minimization of G
- Minimization of 1/G
- Minimization of hybridization function
- High frequency data suppression

$$Diff_i = \sum_{n=0}^M W_n \left| G_{ii}^0(\omega_n) - G_{ii}^{0,exact}(\omega_n) \right|^2$$

- Provide better fitting for small Matsubara frequencies
- Additional weight functions could improve convergence

$$Diff_i = \sum_{n=0}^M W_n \left| \frac{1}{G_{ii}^0(\omega_n)} - \frac{1}{G_{ii}^{0,exact}(\omega_n)} \right|^2$$

- That corresponds to minimization of Hybridization function.
- And for large n

$$\left| G_{ii}^0(\omega_n) - G_{ii}^{0,exact}(\omega_n) \right| \rightarrow \frac{1}{\omega_n^2} \left| \Delta_{ii}(\omega_n) - \Delta_{ii}^{exact}(\omega_n) \right|$$

- Improves asymptotic behavior

$$Diff_i = \sum_{n=0}^M W_n \left| G_{ii}^0(\omega_n) - G_{ii}^{0,exact}(\omega_n) \right|^2$$

- Small Matsubara frequencies if Δ has insulating behavior (min_type=2 in SEET code)

$$W_n = \omega_n$$

- Larger Matsubara frequencies if Δ shows metallic behavior (min_type=1 in SEET code)

$$W_n = 1$$

- Need to perform full optimization for matrix-valued functions
- For small number of bath sites this procedure is unstable
- Better to find diagonal basis