

Nevanlinna Analytic Continuation At Gull's group meeting

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6



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Continuation

NG = -G is a Nevanlinna function ($C^+ \rightarrow C^+$)

Proof. Green's function restricted to C^+ can be formulated by Lehmann representation as,

$$\mathcal{G}(\gamma, z) = \frac{1}{Z} \sum_{m,n} \frac{|\langle m | c_{\gamma}^{\dagger} | n \rangle|^2}{z + E_n - E_m} (e^{-\beta E_n} + e^{-\beta E_m})$$

Denote and notice that, $A = \frac{1}{Z} |\langle m | c_{\gamma}^{\dagger} | n \rangle|^2 (e^{-\beta E_b} + e^{-\beta E_m}) > 0$

Let z = x + yi where y > 0, i.e. $z \in C^+$. Then each summand can be represented as,

$$S = \frac{A}{(x + E_n - E_m) + iy} = \frac{A(x + E_n - E_m - iy)}{(x + E_n - E_m)^2 + y^2}$$
$$Im\{S\} = -\frac{Ay}{(x + E_n - E_m)^2 + y^2} \le 0$$

summing all summands gives $Im\{\mathcal{G}(\gamma, z)\} \leq 0$ and thus $Im\{\mathcal{NG}(\gamma, z)\} \geq 0$ for $z \in \mathcal{C}^+$.

Every Nevanlinna function N admits a representation

$$N(z)=C+Dz+\int_{\mathbb{R}}\left(rac{1}{\lambda-z}-rac{\lambda}{1+\lambda^2}
ight)\mathrm{d}\,\mu(\lambda),\quad z\in\mathcal{H},$$

where C is a real constant, D is a non-negative constant, \mathcal{H} is the upper half-plane

$$G(i\omega_n) = -\frac{1}{\pi} \int \frac{\mathrm{Im}G(\omega)d\omega}{i\omega_n - \omega}$$





Rolf Nevanlinna



Issai Schur



Interpolation Algorithm

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Choose a solution: optimization

0

0

ω

5

0.4

0.2

0.0 0.4

0.2

0.0

5



before optim

after optim

-5

10

(~) 7

0

0.30

(9) W 0.15

0.00

-5

0

0

ω



5

-5

-5

5



FIG. 1. Optimization with 5, 15 and 25 Hardy basis functions. Top panels: resulting spectral functions $A(\omega)$. Bottom panels: real and imaginary part of the exact and fitted parametric functions θ_{M+1} ($\theta_{M+1} : \mathcal{C}^+ \to \overline{\mathcal{D}}$). The needed θ_{M+1} is what would restore our synthetic input







Georg Pick

The Pick criterion: Nevanlinna / Schur interpolants' $\underline{g}(x_i) = y_i \ (x_i \in \mathcal{D}, y_i \in \overline{\mathcal{D}})$ $P_{ij} = \left[\frac{1 - y_i y_j^*}{1 - x_i x_j^*}\right]$

Existence ($P \ge 0$) and Uniqueness (P furthermore singular)



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Noisy data

Continuous fraction interpolation is very sensitive to noise.

If input data does not satisfy Pick criterion, output may or may not be causal.

In practice poles are often close to real axis, evaluation just above the real axis will skip them.

Behavior with noise similar to Padé continued fractions. This is not a method for noisy data.





Hamburger moment problem
Given
$$b = (h_0, h_1, h_2, ...)$$
, construct $\sigma(\omega)$ s.t.
 $h_k = \int \omega^k d\sigma(\omega)$ In our case, $d\sigma(\omega) = A(\omega)d\omega$
 $\mathcal{N}G(i\omega_n) = -\int_0^\beta d\tau G(\tau)e^{i\omega_n\tau}$
 $= -\sum_{k=0}^\infty \frac{(-1)^{k+1}(G^{(k)}(\beta) + G^{(k)}(0))}{(i\omega_n)^{k+1}}$
 $= -\frac{h_0}{i\omega_n} - \frac{h_1}{(i\omega_n)^2} - \frac{h_2}{(i\omega_n)^3} - \cdots$
 $f(z) = -\frac{h_0}{z} - \frac{h_1}{z^2} - \frac{h_2}{z^3} - \cdots - \frac{h_{2N-2}}{z^{2N-1}} - o(\frac{1}{z^{2N-1}})$

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The Carathéodory class of matrix-valued analytic functions in D (or: C^+) is defined as

$$C = \{M(z) : M(z) + M^{\dagger}(z) \ge 0 \quad \forall |z| < 1\} \quad \Leftrightarrow \quad \operatorname{Re}\{x^{\dagger}Mx\} \ge 0$$

 $-iG^{<}(\omega), iG^{>}(\omega)$ are PSD; $iG(z), i\Sigma(z), iM(z)$ (cumulant: defined as $M^{-1}(z) = G^{-1}(z) + F$) are Carathéodory.

Proof.
$$\langle x|-iG^{<}(\omega)|x\rangle = 2\pi \sum_{mn} \frac{e^{-\beta E_n}}{Z} \langle m|\sum_i c_i x_i^*|n\rangle^2 \delta(\omega - E_n + E_m) \ge 0$$

 $\langle x|iG^{>}(\omega)|x\rangle = 2\pi i \sum_{mn} \frac{e^{-\beta E_n}}{Z} \langle n|\sum_i c_i x_i^*|m\rangle^2 \delta(\omega + E_n - E_m) \ge 0$
 $\langle x|i\mathcal{G}(z) + (i\mathcal{G}(z))^*|x\rangle = \frac{1}{Z} \sum_{mn} \frac{2\mathrm{Im}\{z\} (e^{-\beta E_m} + e^{-\beta E_n})}{\mathrm{Im}\{z\}^2 + (\mathrm{Re}\{z\} + E_n - E_m)^2} \langle n|\sum_i c_i x_i^*|m\rangle^2 \ge 0$

$$-iG^{-1}(z) + i(G^{-1}(z))^{\dagger} = -iM^{-1}(z) + iF + i(M^{-1}(z))^{\dagger} - iF^{\dagger}$$
$$= -iM^{-1}(z) + i(M^{-1}(z))^{\dagger}$$

[Gramsch, Potthoff; Phys. Rev. B 92, 235135 (2015)]

Fourier transform of $\delta_C(t, t') \Sigma_{ij}^{\text{HF}}(t)$ Hermitian and z-independent

$$\begin{split} &\langle x|i\Sigma(z) + (i\Sigma(z))^{\dagger}|x\rangle \\ &= \langle x|i\Sigma^{C}(x+yi) - i\Sigma^{C}(x-yi)|x\rangle \\ &= \sum_{s} \frac{2y(\sum_{i} x_{i}h_{is}(0))^{2}}{(h_{ss}-x)^{2}+y^{2}} \geqslant 0 \end{split}$$

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Generalized Schur's algorithm

Assuming $\Psi_i(z_i) = W_i$, $y_i L_i(z) = [I - W_i W_i^{\dagger}]^{-1/2} [\Psi_i(z) - W_i] \cdot [I - W_i^{\dagger} \Psi_i(z)]^{-1} [I - W_i^{\dagger} W_i]^{1/2}$ where $y_i = |z_i|(z_i - z)/(z_i(1 - z_i^* z))$. $L_i(z) \in \mathcal{S}$ by the Schwarz lemma [53]. Define $\Psi_{i+1}(z)$ as $\Psi_{i+1}(z) = [I - K_i K_i^{\dagger}]^{-1/2} [L_i(z) - K_i] \cdot [I - K_i^{\dagger} L_i(z)]^{-1} [I - K_i^{\dagger} K_i]^{1/2}$ where K_i is an arbitrary matrix such that $||K_i|| < 1$.



Carathéodory







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Hubbard dimer (2-site)

 $H_{asym} = H_0 + H_V + H_H + H_{breaksym}$

Here we use the degeneracy-lifted Hamiltonian $H = H_{asym}$ with parameters $\beta = 10$, t = 1, U = 5, $\mu = 0.7$, H = 0.3, $U_a = 0.5$, $\mu_a = 0.2$, $H_a = 0.03$.







FIG. 3. Total spectral function of the Hubbard dimer, obtained in the site basis and in a randomly rotated basis, illustrating the basis independence of the continuation procedure.

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Comments



- This work eliminates the leading problem in analyzing diagrammatic / stat mech / finite temperature calculations with semi-analytic formalisms (GW, FLEX, T-Matrix, fRG, etc)
- Noise case was untouched.

Ideas from Nevanlinna?

holographic mapping, hardy space and Fourier coefficients [Ying 22; 2202.09719];

Project to the causal space then choose a solution

(1) handling Schur interpolant, systematic way to project / eliminate noise [control theory / operator theory]

(2) projecting Matsubara data to the Nevanlinna space (may be NP hard and useless: dense causal volume in the highdimensional Matsubara data space; many causal data are 'close' enough to the given noisy data but give unmovable bad spectra)

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How to run Nevanlinna



Dependencies: C++, CMake; Dakota (opt) Linked libraries: MPFR, GMPXX, GMP; FFTW3 (opt)

- 1. Compile the program inside "Nevanlinna_Schur". Check line 74 of nevanlinna.h and line 28,29 of sobolev.cpp, making sure N_real_=N and omega_min=omega_max=length. \lambda could be adjusted in line 113 of sobolev.cpp.
- 2. Run "./Nevanlinna < input.txt", the fourth "coeff" in "input.txt" is the name of the coefficient file that will be used in Hardy optimization.
- 3. Put the generated "coeff", "Nev_opt.py" and "Nev_opt_result.py" under the same directory used for optimization.
- 4. Run "Nev_opt.py". *It* creates folder "dakota"; writes three files "dakota/<u>dakota_pstudy.in</u>" [optimization method, number of parameters which is 4*basis -> Re/Im of ak and bk, initial values], "dakota/driver" [it runs the driver program "sobolev.cpp" inside the iteration folder "dakota/0", "dakota/1", ...] and "dakota/params.template" [the template for dakota to fill in trial parameters in each iteration]; and runs "jobopt".
- 5. After optimization, run "Nev_opt_result.py". It extracts result files and puts them into "opt_result" folder.
- 6. Run the "job" file created **by** "Nev_opt_result.py". This would finally give "Aopt.txt" in the same directory, which is the optimized spectrum.

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